

# Load-dependent destabilization of the $\gamma$ -rotor shaft in $F_0F_1$ ATP synthase revealed by hydrogen/deuterium-exchange mass spectrometry

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$F_0F_1$  is a membrane-bound molecular motor that uses proton-motive force (PMF) to drive the synthesis of ATP from ADP and  $P_i$ . Reverse operation generates PMF via ATP hydrolysis. Catalysis in either direction involves rotation of the  $\gamma$  shaft that connects the  $\alpha_3\beta_3$  head and the membrane-anchored  $c_r$  ring. X-ray crystallography and other techniques have provided insights into the structure and function of  $F_0F_1$  subcomplexes. However, interrogating the conformational dynamics of intact membrane-bound  $F_0F_1$  during rotational catalysis has proven to be difficult. Here, we use hydrogen/deuterium exchange mass spectrometry to probe the inner workings of  $F_0F_1$  in its natural membrane-bound state. A pronounced destabilization of the  $\gamma$  C-terminal helix during hydrolysis-driven rotation was observed. This behavior is attributed to torsional stress in  $\gamma$ , arising from  $\gamma$ - $\alpha_3\beta_3$  interactions that cause resistance during  $\gamma$  rotation within the apical bearing. Intriguingly, we find that destabilization of  $\gamma$  occurs only when  $F_0F_1$  operates against a PMF-induced torque; the effect disappears when PMF is eliminated by an uncoupler. This behavior resembles the properties of automotive engines, where bearings inflict greater forces on the crankshaft when operated under load than during idling.

molecular motor | molecular bearing | conformational dynamics | destabilized helix | rotational resistance

Bacteria, mitochondria, and chloroplasts share a similar  $F_0F_1$ -ATP synthase architecture (1–3) (Fig. 1A). The *Escherichia coli* enzyme considered in this work has the rotor composition  $\gamma\epsilon c_{10}$  (4). The stator consists of the  $\alpha_3\beta_3$  catalytic head, the  $\delta b_2$  peripheral stalk, and the  $a$  subunit. The  $\gamma\epsilon$ -rotor shaft connects  $\alpha_3\beta_3$  with the  $ac_{10}$  proton translocator (3, 5). The intact molecular motor has a mass of 525 kDa.  $F_0F_1$  investigations are typically conducted under hydrolysis conditions (1, 6, 7), where  $\gamma\epsilon c_{10}$  rotation is driven by the  $\beta$  subunits that cycle through a series of conformations (8). ATP hydrolysis triggers power strokes that involve consecutive interactions of the  $\beta$ -levers with  $\gamma$  (3, 6, 9).

X-ray crystallography (5, 8, 10, 11) and cryoelectron microscopy (12, 13) have provided valuable snapshots of subcomplexes and solubilized  $F_0F_1$ . Molecular dynamics (MD) simulations (14–16) can help interpret such data in a dynamic context. However, there is scant information on how the conformation of  $F_0F_1$  is affected by internal and external forces while working under physiological conditions, where ATP/ADP interconversion is coupled with transmembrane proton transport in the presence of proton-motive force (PMF). With few exceptions (1, 3, 17), previous experiments on the rotary mechanism have focused on isolated  $F_1$ .

Of particular interest is the question of how moving protein surfaces within molecular motors interact with each other. The irregular topologies and specific binding interactions of biomolecular interfaces complicate a rigorous discussion of dissipative effects in terms of classical friction models (18). It seems likely that sliding motions at such interfaces will be associated with resistance as individual residues clash with one another, and as transient binding interactions are formed and disrupted (19). Rotation of  $\gamma$  within the  $\alpha_3\beta_3$  apical bearing of  $F_0F_1$  represents a prime example of such a

scenario. Along these lines, recent MD simulations provided evidence for friction during torque-driven rotation of  $\gamma$  (15). More generally, both solvent friction and internal friction have been discussed in the context of conformational changes for monomeric proteins (19–21). For addressing these issues in the context of  $F_0F_1$ , it is necessary to interrogate the structure and dynamics of the molecular motor in situ while maintaining the system under different catalytically active and inhibited conditions.

Hydrogen/deuterium exchange (HDX) mass spectrometry (MS) is a sensitive approach for probing changes in protein structure and dynamics in response to external stimuli (22, 23). This technique monitors backbone deuteration in the presence of  $D_2O$ . Rigid protein segments exchange more slowly than regions that are not as structurally stable (24). Under continuous labeling conditions, the deuteration kinetics can be monitored in a spatially resolved fashion by subjecting aliquots to peptic digestion, followed by liquid chromatography and MS analysis of the resulting peptides. The use of HDX/MS is well established for protein binding experiments and several other types of investigations (22, 23); however, this technique is underused for mechanistic studies on molecular machines (25). Ryrie and Jagendorf (26) monitored hydrogen/tritium exchange in chloroplast  $F_0F_1$  four decades ago, but those experiments did not yield spatially resolved information.

Here, we use HDX/MS to compare the conformational dynamics of  $F_0F_1$  under various catalytically active and inhibited conditions. The experiments were conducted using vesicles that represent the natural *E. coli* membrane environment. Our data

## Significance

$F_0F_1$ , or ATP synthase, is often referred to as the “world’s smallest motor.” Similar to automotive engines, it employs a rotating shaft that interacts with mechanical actuators. When operating a combustion engine under load, the bearings exert significant forces on the crankshaft, leading to enhanced mechanical stress. Here, we demonstrate that analogous load-dependent effects occur in molecular motors. When  $F_0F_1$  pumps protons against a transmembrane gradient, the rotor shaft undergoes structural destabilization attributed to resistive forces in its apical bearing. The effect disappears when the transmembrane gradient opposing proton pumping is short-circuited by an uncoupler, as predicted by fundamental principles of mechanics. Our observations highlight fascinating parallels between engine operation on the macroscale and the nanoscale.

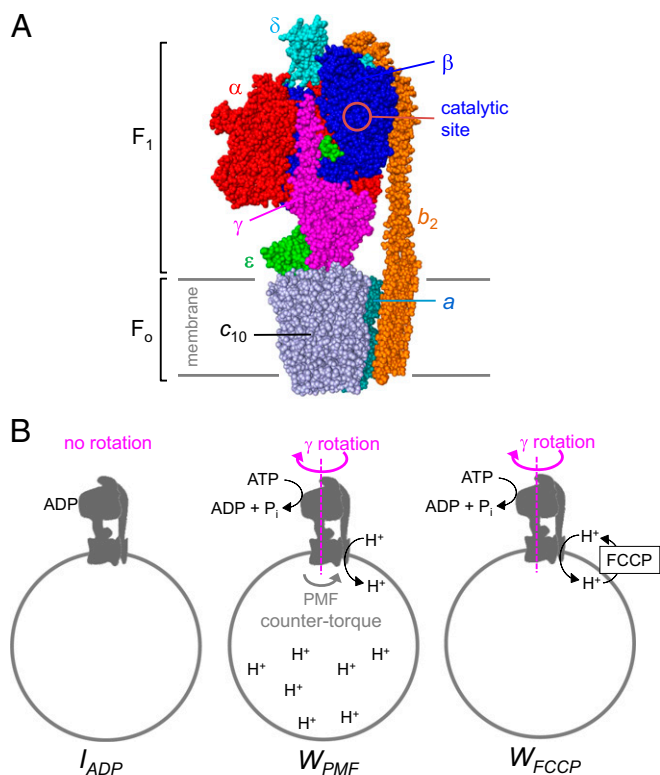
Author contributions: S.V., S.D.D., and L.K. designed research; S.V., Y.B., and S.D.D. performed research; Y.B. and S.D.D. prepared and characterized  $F_0F_1$  membrane vesicles; S.V. and L.K. analyzed data; and L.K. wrote the paper.

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**Fig. 1.** Subunit architecture of  $F_0F_1$  ATP synthase from *E. coli*. (A) Composite model assembled from subcomplex structures [Protein Data Bank (PDB) ID codes 3OAA, 3J0J, 1C17, 2XOK, and 2WSS]. One pair of  $\alpha\beta$  subunits facing the observer was removed to illustrate how  $\gamma$  extends into the apical bearing formed by the  $\alpha_3\beta_3$  head. During rotational catalysis,  $\epsilon$  (partially obstructed by  $\gamma$ ) is likely folded down toward  $c_{10}$ , rather than being extended as shown here (10). (B) Cartoon of membrane vesicle-bound  $F_0F_1$  under different experimental conditions.  $I_{ADP}$  represents ADP-inhibited  $F_0F_1$ ;  $W_{PMF}$  represents catalytically active  $F_0F_1$  that pumps protons against a PMF-mediated counter-torque; and  $W_{FCCP}$  refers to catalytically active  $F_0F_1$  that pumps protons in the presence of the uncoupler FCCP, which prevents PMF buildup.

reveal that the  $\gamma$  C-terminal helix experiences load-dependent destabilization. This effect reveals that rotation of the helix within the  $\alpha_3\beta_3$  apical bearing is hindered by interchain contacts. We believe this investigation to be the first in situ HDX/MS investigation of a catalytically active rotational molecular motor.

## Results and Discussion

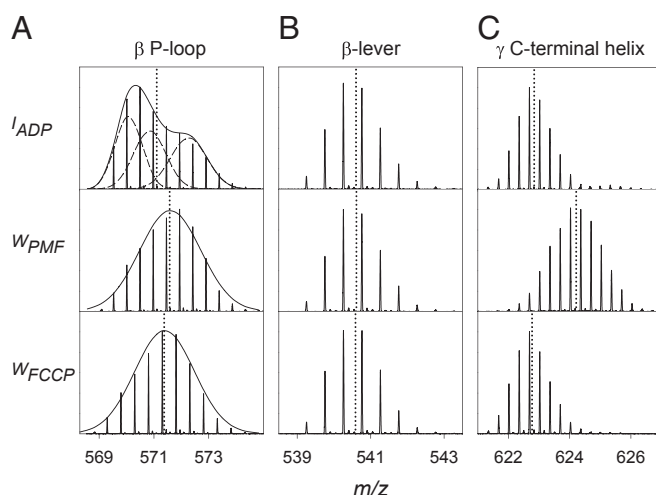
**HDX/MS of  $F_0F_1$  in Membrane Vesicles.** Deuteration measurements were conducted on  $F_0F_1$  in inside-out *E. coli* membrane vesicles, with focus on three conditions (Fig. 1B): (i) The presence of ADP (without ATP) produces the inhibited state  $I_{ADP}$ , where Mg-ADP and azide remain permanently bound in at least one catalytic site (27); (ii) the  $W_{PMF}$  state is catalytically active (“working”; i.e., protons are pumped into the vesicle against PMF); and (iii)  $W_{FCCP}$  also represents catalytically active  $F_0F_1$ , but PMF buildup is prevented by the uncoupler carbonyl cyanide-4-(trifluoromethoxy) phenylhydrazone (FCCP). An ATP regeneration system was used to ensure that  $W_{PMF}$  and  $W_{FCCP}$  underwent ATP hydrolysis with  $k_{cat} = 11 \pm 1 \text{ s}^{-1}$ , resulting in  $\sim 30,000$  catalytic events per enzyme during the 45-min HDX time window. This value corresponds to  $\sim 10,000$  complete  $\gamma$  rotations for both  $W_{PMF}$  and  $W_{FCCP}$ , whereas rotation is blocked under  $I_{ADP}$  conditions.

Control measurements were extended to two additional types of samples: adenosine 5'-( $\beta$ - $\gamma$ -imido)triphosphate (AMP-PNP)-inhibited  $F_0F_1$ , as well as Mg<sup>2+</sup>-depleted preparations (SI Appendix, Fig. S1). These control measurements confirmed that HDX/MS is capable

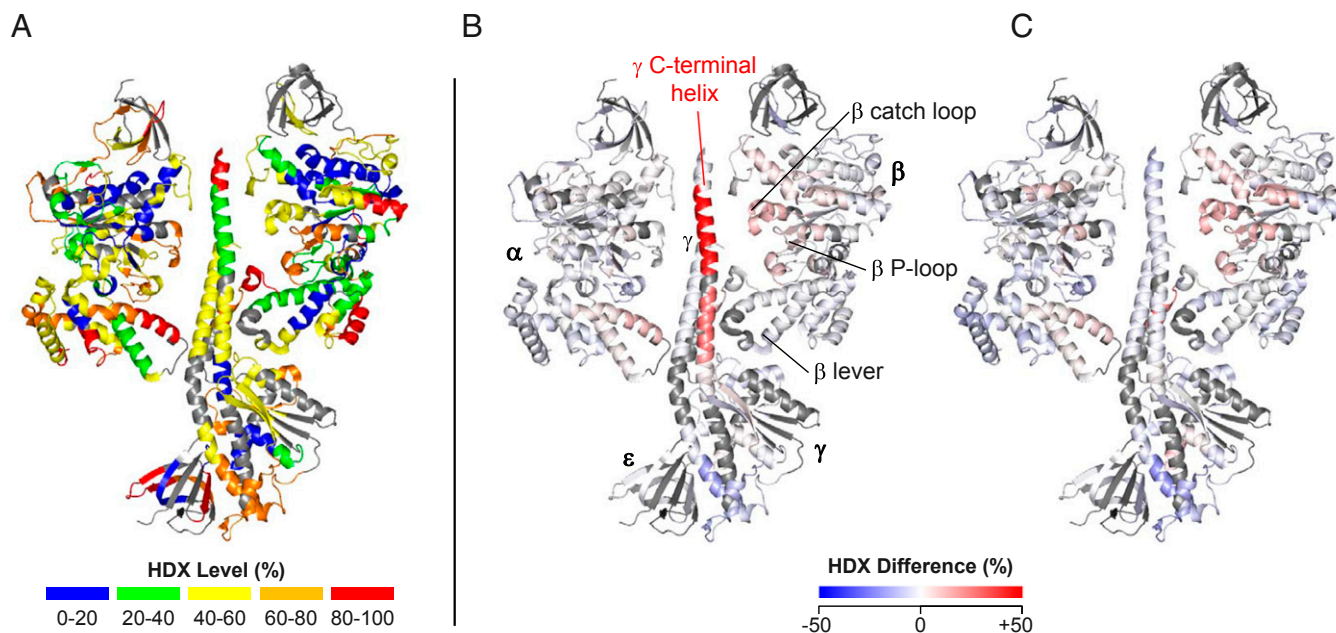
of pinpointing changes in structure and dynamics of vesicle-bound  $F_0F_1$  with high fidelity (SI Appendix, Fig. S1A).

The use of natural *E. coli* membranes resulted in peptic digests containing a large number of peptides from proteins other than  $F_0F_1$ . Ribosomal proteins, in particular, were present in high abundance, despite the use of sucrose gradient ultracentrifugation during membrane isolation (details are provided in SI Appendix, Methods). A multidimensional analysis work flow using peptide separation based on  $m/z$ , retention time, and ion mobility drift time was required to cope with these highly complex samples (SI Appendix, Fig. S2). These digestion and analysis conditions consistently yielded a total of 203  $F_0F_1$  peptides with signal-to-noise ratios that were adequate for providing highly reproducible HDX/MS data. Sequence coverage for the extramembrane subunits of  $F_0F_1$  was high ( $\alpha$ , 83%;  $\beta$ , 81%;  $\gamma$ , 74%;  $\delta$ , 77%;  $\epsilon$ , 48%;  $b$ , 58%). Unfortunately, only a few peptides were detected for membrane-embedded subunits, such that a meaningful characterization of  $a$  and  $c$  was not possible. Low digestion yields for transmembrane segments are common in HDX/MS (28). The use of natural membranes in this work is particularly challenging (29). Notably, most earlier membrane protein HDX/MS investigations used detergent-solubilized samples (30) or purified membrane surrogates, such as nanodiscs (31).

**Deuteration Behavior of Selected Peptides.** It is instructive to look initially at selected unprocessed HDX/MS data. The P-loop peptide  $\beta^{148}\text{FGGAGVGKTVNM}^{159}$  is involved in nucleotide binding (8). In the  $I_{ADP}$  state, this region displays asymmetrical HDX distributions (Fig. 2A) that are consistent with the three stable catalytic site conformations  $\beta_{right}$ ,  $\beta_{loose}$ , and  $\beta_{open}$  (32). We attribute the most extensively deuterated P-loop component to  $\beta_{loose}$ , which is known to adopt a distorted conformation with several disrupted H-bonds (8). Rotational averaging under  $W_{PMF}$  and  $W_{FCCP}$  conditions causes coalescence into a unimodal HDX envelope (Fig. 2A). Similar effects were observed for two other peptides close to the catalytic sites:  $\beta^{177}\text{AGVGERTREGNDF}^{189}$  and  $\alpha^{360}\text{FNAGIRPAVNPGIS}^{373}$ . Regions that are insensitive to changes in conditions include the  $\beta$ -levers (Fig. 2B). The  $\gamma$  C-terminal helix shows greatly enhanced HDX for  $W_{PMF}$ , but not for  $W_{FCCP}$  or  $I_{ADP}$  (Fig. 2C). Results for  $\epsilon$  are illustrated in SI Appendix, Fig. S3.



**Fig. 2.** HDX mass spectra of selected peptides in the  $I_{ADP}$ ,  $W_{PMF}$ , and  $W_{FCCP}$  states. (A) Data for a peptide comprising the active site P-loop ( $\beta^{148}\text{FGGAGVGKTVNM}^{159}$ ) after 1.5 min of deuteration. Gaussian deconvolution of the  $I_{ADP}$  spectrum reveals the presence of three equally populated noninterconverting conformers, attributed to  $\beta_{right}$ ,  $\beta_{loose}$ , and  $\beta_{open}$ . (B)  $\beta$ -Lever region ( $\beta^{380}\text{DELSEEDK}^{388}$ ) at  $t = 45$  min. (C)  $\gamma$  C-terminal region ( $\gamma^{260}\text{LQLVYINKARQASITQE}^{275}$ ) at  $t = 45$  min. Vertical dotted lines represent centroid  $m/z$  values.



**Fig. 3.** (A) HDX levels of  $I_{ADP}$  for an HDX period of 45 min. Colors represent deuteration percentages as indicated in the legend; these data were derived from the centroids of the corresponding isotope distributions. To simplify the graphic representation, this figure only displays HDX/MS data for one  $\alpha$  subunit and one  $\beta$  subunit, as well as  $\gamma$  and  $\epsilon$  (PDB ID code 3OAA). (B) Deuterium difference map of  $W_{PMF}$  vs.  $I_{ADP}$ . (C) Deuterium difference map of  $W_{FCCP}$  vs.  $I_{ADP}$ . Dark red coloring of the  $\gamma$  C-terminal helix in B highlights dramatically enhanced deuteration of  $\gamma^{260}$ LQLVYNKARQASITQEL $^{276}$  under  $W_{PMF}$  compared with  $I_{ADP}$  conditions. Regions not covered by peptide mapping are shown in dark gray.

**Comprehensive Overview of HDX Patterns.** For discussing the HDX properties of  $F_0F_1$ , it is convenient to adopt  $I_{ADP}$  as a reference state (Fig. 3A), yielding the difference maps of Fig. 3B and C.  $W_{PMF}$  conditions are known to induce elastic deformation of the  $\gamma$  globular bottom domain (33), because the  $\beta$ -lever action is opposed by a  $c_{10}$  countertorque (3). For  $W_{FCCP}$ , this mechanical stress is greatly reduced because the  $c_{10}$  countertorque is eliminated. These considerations suggest that  $F_0F_1$  power transmission elements might exhibit PMF-dependent HDX characteristics. Unexpectedly, Fig. 3B and C reveals that the HDX behavior of key power transmission elements [ $\gamma$  bottom domain and  $\beta$ -levers (3)] is largely identical with and without PMF (Fig. 3B and C). We conclude that much of the  $F_0F_1$  H-bonding network resists perturbation by intramolecular forces encountered during rotational catalysis.

Interestingly, major PMF-dependent HDX changes are seen for parts of the  $\gamma$  C-terminal helix, with deuteration levels that are  $\sim 40\%$  higher under  $W_{PMF}$  conditions than in the  $W_{FCCP}$  and  $I_{ADP}$  states (Figs. 3B and C and 4). This effect is most pronounced for the range  $\gamma^{260}$ LQLVYNKARQASITQEL $^{276}$ , which appears in dark red in Fig. 3B. The corresponding region is covered by five partially overlapping peptides (SI Appendix, Fig. S5). The greatly enhanced deuteration of this  $\gamma$  region reveals the occurrence of structural perturbations when rotational catalysis proceeds in the presence of PMF (Fig. 3B). No such destabilization is observed when the rotor shaft is stationary ( $I_{ADP}$ ) or when rotation takes place in the absence of PMF ( $W_{FCCP}$ ; Fig. 3C).

**Characterizing the Structure of PMF-Destabilized  $\gamma$ .** For exploring the nature of PMF-induced structural perturbations in the  $\gamma$  C-terminal helix, it is helpful to consider the deuteration behavior of  $\gamma$  in the context of classical HDX theory (24, 34, 35). Backbone NH groups in regions that are permanently disordered are known to undergo rapid HDX with an overall rate constant  $k_{HDX}$  that approaches the “chemical” rate constant of  $k_{ch} \approx 30 \text{ s}^{-1}$  at pD 8 (35). Such rapid deuteration causes a burst phase in continuous labeling experiments. This behavior is displayed by the extreme C terminus of  $\gamma$  ( $^{279}$ IVSGAAAV $^{286}$ ; Fig. 4B), which shows

a burst amplitude of  $\sim 70\%$ . In other words,  $\gamma^{279}$ IVSGAAAV $^{286}$  is largely disordered in solution, regardless of experimental conditions.

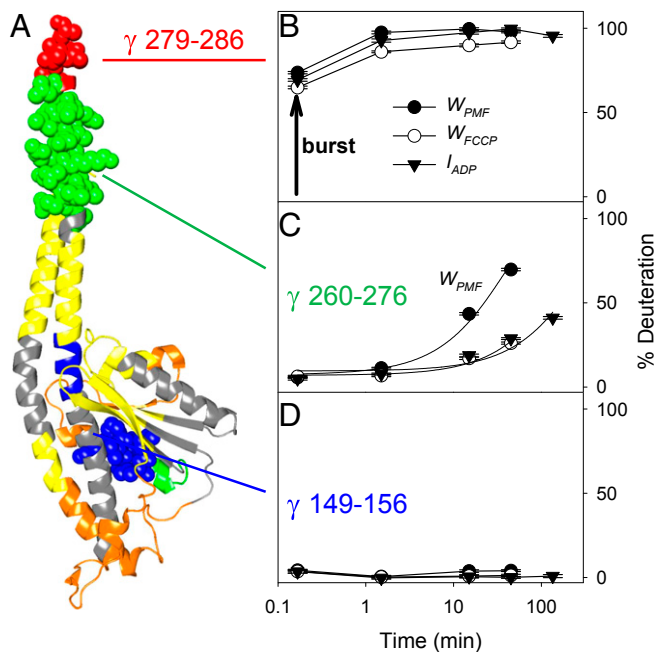
Deuteration of amide NH groups in folded regions is mediated by transient H-bond opening/closing fluctuations, as envisioned by the native state mechanism (24, 35):



where  $k_{cl} \gg k_{op}$ . In the EX2 regime encountered here, the resulting overall rate constant is  $k_{HDX} = (k_{op}/k_{cl}) \times k_{ch}$  and the free energy associated with H-bond opening is  $\Delta G^\circ = -RT \ln(k_{HDX}/k_{ch})$ . Amides that follow this native state mechanism do not show a burst phase (i.e., deuteration for short HDX times is close to zero). For extremely stable regions,  $k_{op}/k_{cl}$  is so small that deuteration remains virtually undetectable throughout the entire experimental time window. Such a case is encountered for  $\gamma^{149}$ IGPVKVML $^{156}$ , which is part of the globular  $\gamma$  bottom domain (Fig. 4D).

The  $\gamma$  C-terminal helix region that undergoes PMF-dependent destabilization ( $\gamma^{260}$ LQLVYNKARQASITQEL $^{276}$ ) shows deuteration kinetics in-between these two extremes (Fig. 4C). None of the profiles in Fig. 4C exhibits a burst phase, implying that residues  $\gamma$  260–276 remain predominantly folded, with occasional H-bond opening/closing transitions. These transient opening events are much more extensive under  $W_{PMF}$  conditions than for  $W_{FCCP}$  and  $I_{ADP}$ . The heavy extent of deuteration and greatly increased HDX rate of  $W_{PMF}$  are readily apparent in Fig. 4C. The  $k_{HDX}$  values estimated from the exponential fits of Fig. 4C imply that hydrogen bonds in this region get destabilized by  $\sim 4 \text{ kJ}\cdot\text{mol}^{-1}$  ( $1\text{--}2 k_B T$ ) on average, from  $\Delta G^\circ \approx 32 \text{ kJ}\cdot\text{mol}^{-1}$  in the  $W_{FCCP}$  state to  $28 \text{ kJ}\cdot\text{mol}^{-1}$  per hydrogen bond under  $W_{PMF}$  conditions.

**Evidence for Hindered Rotation of  $\gamma$ .** What causes the destabilization of the  $\gamma$  C-terminal helix? Any mechanistic explanation has to address the observation that rotation, per se, is insufficient for causing this destabilization ( $W_{FCCP}$ ; Fig. 3C). Instead, destabilization is observed only when  $F_0F_1$  operates against a PMF-mediated countertorque



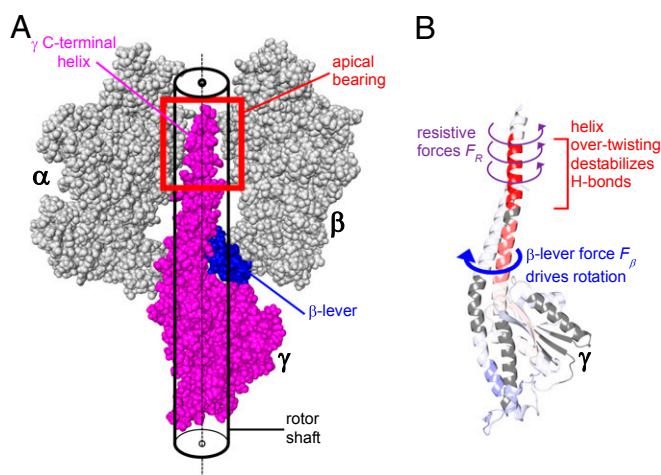
**Fig. 4.** Deuteration behavior of selected  $\gamma$  segments. (A) X-ray structure of  $\gamma$ , colored as in Fig. 3A. Three peptides are highlighted in space-fill representation:  $\gamma^{279}$ IVSGAAV $^{286}$ ,  $\gamma^{260}$ LQLVYNKARQASITQEL $^{276}$ , and  $\gamma^{149}$ IGPVKVM $^{156}$ . (B–D) HDX kinetics of these peptides are shown for  $I_{ADP}$ ,  $W_{PMF}$ , and  $W_{FCCP}$  conditions. Lines in C are single-exponential fits, subject to the constraint that the deuteration level for  $t = \infty$  is 100%. The corresponding rate constants are  $6 \times 10^{-5} \text{ s}^{-1}$ ,  $46 \times 10^{-5} \text{ s}^{-1}$ , and  $9 \times 10^{-5} \text{ s}^{-1}$ , respectively. The vertical arrow in B represents burst phase labeling with an amplitude of  $\sim 70\%$ .

( $W_{PMF}$ ; Fig. 3B). Different reasons may be considered when trying to account for this behavior:

- i) Direct power transmission: The most strongly destabilized region (Fig. 3B, dark red) is not directly involved in  $\beta/\gamma$  power transmission (3, 36, 37). Even  $F_0F_1$  constructs with severely truncated  $\gamma$  C termini can still generate significant torque (36, 38, 39). Those previous findings imply that the PMF-induced destabilization of the  $\gamma$  C-terminal helix seen in our experiments is not a direct power transmission effect.
- ii) Reverse rotation during dwell periods: Individual power strokes cause  $\gamma$  to turn  $120^\circ$  in a clockwise direction, when viewed from the top of the  $\alpha_3\beta_3$  crown (40). Each power stroke is followed by a dwell, during which none of the  $\beta$ -levers actively drives  $\gamma$  rotation (6). One may contemplate whether PMF can cause reverse (counterclockwise) rotation of  $c_{10}$  by a small angle ( $\ll 120^\circ$ ) during these dwells, as suggested by some early experiments (41). Could such reverse rotation events destabilize the  $\gamma$  C-terminal helix? The  $\gamma c_{10}$  rotor is remarkably compliant (“soft”) in the globular bottom domain of  $\gamma$ , as required for elastic power transmission (3). Small reverse rotation events would likely be absorbed in this compliant region, such that their effects would not be felt at the opposite end of  $\gamma$ . Thus, it appears implausible that reverse rotation events could be responsible for the PMF-dependent destabilization of the  $\gamma$  C-terminal helix.
- iii) Permanent stalling of the  $\gamma$  C terminus: Hilbers et al. (37) demonstrated that rotational catalysis proceeds unimpeded in  $F_1$  constructs that had the  $\gamma$  C terminus cross-linked with  $\alpha$  via an engineered disulfide bridge. In those constructs, the torque provided by  $\alpha_3\beta_3$  permanently unfolds parts of the  $\gamma$  C-terminal helix, causing  $\gamma$  to undergo  $\phi/\psi$  swivel motions instead of rotating as a whole (37). One has to consider the possibility that comparable effects are encountered for wild-type  $F_0F_1$  in our

$W_{PMF}$  experiments, (i.e., complete stalling of the  $\gamma$  C terminus and unraveling of the  $\gamma$  C-terminal helix). Catalysis in the cross-linked constructs of Hilbers et al. (37) proceeds with permanent opening of approximately eight H-bonds in the  $\gamma$  C-terminal region. Such a situation would cause an HDX burst phase. The absence of a burst in Fig. 4C argues against permanent unfolding, making it unlikely that PMF induces complete stalling of the  $\gamma$  C terminus in our experiments. The work of Hilbers et al. (37) nonetheless provides important clues for the current discussion. Specifically, one can consider a scenario where  $\gamma$  is not permanently stalled, but experiences hindrance as the C-terminal helix rotates within the  $\alpha_3\beta_3$  apical bearing. In the following section, we make the case that such a model is consistent with our data.

- iv) Rotational resistance in the apical bearing: The movement of closely spaced protein surfaces relative to each other is accompanied by side-chain clashes and transient binding/dissociation events (19), likely giving rise to intermittent motion in a stick-slip regime (18). MD simulations indicate that such friction-like effects are also encountered during rotation of the  $\gamma$  C-terminal helix within the  $\alpha_3\beta_3$  apical bearing (15). The simulated  $\gamma$  rotation in the study by Okazaki and Hummer (15) is orders of magnitude faster than in our experiments, thereby complicating comparisons of the data obtained. Nonetheless, our findings support the basic conclusion of Okazaki and Hummer (15) that rotation of  $\gamma$  is associated with mechanical resistance. The  $F_0F_1$  apical bearing region is highlighted in Fig. 5A; it includes the sleeve surrounding the extreme  $\gamma$  C terminus (8), as well as adjacent elements, such as the  $\beta$ -catch loops (42, 43). Rotation will be accompanied by steric clashes as  $\gamma$  side chains are forced past the side chains of  $\alpha_3\beta_3$  within the tight annular gap of the bearing (15). Intermolecular H-bonds and other noncovalent  $\gamma \cdots \alpha_3\beta_3$  linkages (43, 44) may amplify these resistive effects. Fig. 5B illustrates how the resulting resistive forces ( $F_R$ , purple) will hinder rotation of  $\gamma$ . These  $F_R$  oppose the  $\beta$ -lever force ( $F_\beta$ , blue), thereby causing overtwisting of the  $\gamma$  C-terminal helix. The torsional stress generated under these conditions will destabilize backbone H-bonds such that HDX rates are enhanced (Fig. 5B). Protein-protein contacts will also hinder



**Fig. 5.** (A) Close-up view of  $\gamma/\alpha_3\beta_3$  contacts (PDB ID code 3OAA after hydrogen addition; only one  $\alpha\beta$  pair is shown). Surface roughness in the apical bearing causes  $F_R$  during  $\gamma$  rotation. (B) Interplay between  $F_R$  exerted by the apical bearing and  $F_\beta$  induces torsional stress in the  $\gamma$  C-terminal helix during each power stroke. We propose that this torsional stress destabilizes H-bonds, and thereby accelerates HDX in this region. The coloring of  $\gamma$  in B is consistent with the coloring of  $\gamma$  in Fig. 3B.

rotation elsewhere along the shaft, but the apical region is most vulnerable because the apical region is where  $\gamma$  tapers from a coiled coil into a single helix.

Scenario *iv* where  $\gamma$  encounters rotational resistance can explain the observed PMF dependence. An eccentric force acting on a shaft results in a stable rotation axis only if the shaft is supported by suitable bearings (45). This general principle applies equally to piston-driven crankshaft rotation and to  $\beta$ -lever-driven  $\gamma$  rotation in  $F_0F_1$ . Combustion engine measurements revealed that forces exerted by bearings on the crankshaft are much greater under load than during idling (45). Surface protrusions on either side of the bearing/shaft interface are more likely to interact with each other in the presence of a load, where the pressure inside the bearing is significantly enhanced (45, 46). In extreme cases, these interactions can induce material degradation (46). Our HDX data suggest that analogous effects are experienced by the  $\gamma$ -rotor shaft. The apical  $\alpha_3\beta_3$  bearing will exert greater mechanical pressure on the  $\gamma$  C-terminal helix in the presence of a PMF load. This bearing pressure will promote interlocking of side chains in the annular gap, thereby enhancing  $F_R$  (purple arrows in Fig. 5B). As indicated in Fig. 5B, the resulting interplay of  $F_R$  and  $F_\beta$  causes overtwisting of the  $\gamma$  C-terminal helix during each power stroke, thereby destabilizing H-bonds and increasing HDX rates. Bearing pressure is lower in the absence of PMF, such that contacts at the  $\gamma/\alpha_3\beta_3$  interface become less tight. The resulting drop in  $F_R$  diminishes the extent of helix overtwisting during power strokes, thereby lowering HDX levels under  $W_{FCCP}$  conditions. A semiquantitative framework that further dissects the interplay between PMF and  $\gamma$  overtwisting is outlined in *SI Appendix, Fig. S6*.

In summary, the PMF-induced destabilization of the  $\gamma$  C-terminal helix is a nontrivial phenomenon. Of all the explanations considered above, scenario *iv* is the most plausible one. It attributes the observed destabilization to rotational resistance, consistent with crystallographically detected contacts between  $\gamma$  and  $\alpha_3\beta_3$  (43, 44), and with nonspecific friction seen in MD simulations (15).

## Conclusions

$F_0F_1$ -ATP synthase in a cellular environment normally operates in the presence of PMF. Thus, the findings of this work imply that a certain degree of rotational resistance is intrinsic to  $F_0F_1$  operation *in vivo*. Our experiments could only explore ATP hydrolysis-driven rotation, but resistance will likely also be encountered during ATP synthesis.

The HDX kinetics indicate that  $\sim 16$  hydrogen bonds in the  $\gamma$  C-terminal region of  $F_0F_1$  become destabilized by  $\sim 4$  kJ $\cdot$ mol $^{-1}$  under  $W_{PMF}$  conditions, corresponding to an overall free energy “penalty” of  $\sim 64$  kJ $\cdot$ mol $^{-1}$ , which is roughly equivalent to the hydrolysis of two ATP molecules (3). Evidently, it is impossible that this amount of free energy is invested during each single ATP hydrolysis event, or during each 360° rotation of  $\gamma$ . Instead, we propose that this 64 kJ $\cdot$ mol $^{-1}$  destabilization reflects accumulated torsional stress in  $\gamma$  that builds up gradually and that persists over extended time periods while  $F_0F_1$  is catalytically active. In other words, the rotational resistance encountered under  $W_{PMF}$  conditions prevents  $\gamma$  from returning to a torsionally relaxed conformation between power strokes. This scenario is consistent with the well-established view that torsional elasticity

allows  $\gamma$  to serve as an energy reservoir (3, 33). The functional significance of strong interactions in the apical bearing is revealed by the findings that disruption of contacts in this region by either truncation of  $\gamma$  or specific amino acid changes to  $\beta$  or  $\gamma$  impairs or prevents ATP-dependent H $^+$ -pumping in membrane vesicles and oxidative phosphorylation *in vivo* (42, 43, 47).

Experiments on isolated  $F_1$  suggest a catalytic efficiency close to 100% (3, 48). Other data indicate that the efficiency is somewhat lower (6), leaving room for dissipative phenomena. Our proposal of persisting torsional stress in the  $\gamma$  C-terminal helix is consistent with highly efficient operation of  $F_0F_1$ . The total free energy converted by each  $F_0F_1$  enzyme in our experiments (with 30,000 turnover events) is  $\sim 10^6$  kJ $\cdot$ mol $^{-1}$ . Thus, the percentage of free energy associated with torsional stress accumulation is negligibly small (on the order of 64 kJ $\cdot$ mol $^{-1}/10^6$  kJ $\cdot$ mol $^{-1}$ ).

Wild-type  $F_0F_1$  under  $W_{PMF}$  conditions shares interesting parallels with the disulfide constructs used by Hilbers et al. (37), although the extent of the effects encountered in either case is different: (i)  $W_{PMF}$  experiences mechanical resistance during rotation of the  $\gamma$  C-terminal helix, whereas the disulfide constructs operate with a permanently stalled  $\gamma$  C terminus; (ii) rotational obstacles in  $W_{PMF}$  give rise to structural destabilization of the  $\gamma$  C-terminal helix, whereas the disulfide constructs experience complete unfolding of a large helix segment; and (iii)  $W_{PMF}$  conditions carry a moderate free energy penalty of  $\sim 64$  kJ $\cdot$ mol $^{-1}$ , but the persisting nature of the structural destabilization allows wild-type  $F_0F_1$  to perform highly efficient catalysis. Unraveling of the helix in the disulfide constructs carries a very large free energy penalty that does not interfere with catalysis due to the persisting nature of the structural change (the helix does not refold after the initial unraveling event). Overall, the behavior displayed by the permanently stalled constructs of Hilbers et al. (37) bolsters the validity of our conclusions; their study demonstrates that efficient catalysis is possible even under conditions that are much more severe than the gentle rotational resistance deduced from our HDX/MS data for wild-type  $W_{PMF}$ .

The load-dependent rotor destabilization seen here for  $\gamma$  represents a prototypical power train feature (45, 46). Our findings highlight the fact that rotor bearings in macroscopic engines and molecular motors share common operational features. HDX/MS is well suited for interrogating these and other aspects of protein-based nanomachines. In the future, it will be interesting to apply this approach to other types of molecular motors, including flagellar systems.

## Methods

The deuteration kinetics of  $F_0F_1$  were monitored under various inhibited (*I*) and working (*W*) conditions [ $I_{ADP}$ ,  $W_{PMF}$ ,  $W_{FCCP}$ ,  $I_{AMP-PNP}$ , and the Mg $^{2+}$ -depleted inhibited state ( $I_{Mg-dep}$ )]. Inside-out *E. coli* membrane vesicles were exposed to 90% D $_2$ O-based labeling buffer. Aliquots were removed at various time points. Peptic digestion was carried out offline under acid-quench conditions, and peptides were separated using reverse-phase chromatography on a nano-Acquity UPLC system (Waters). Ion mobility separation and mass analysis were conducted on a Synapt G2 electrospray quadrupole time-of-flight mass spectrometer (Waters). Details regarding sample preparation and experimental methods are provided in *SI Appendix*.

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